

# Convergence of Filtered Spherical Harmonic Equations for Radiation Transport

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# Outline & References

- Filtered  $P_N$  equations<sup>1</sup>
- Convergence analysis
  - Modified equation<sup>2</sup>
  - Galerkin estimate<sup>3</sup>
  - Convergence estimates
- Numerical experiments using StaRMAP<sup>4</sup>

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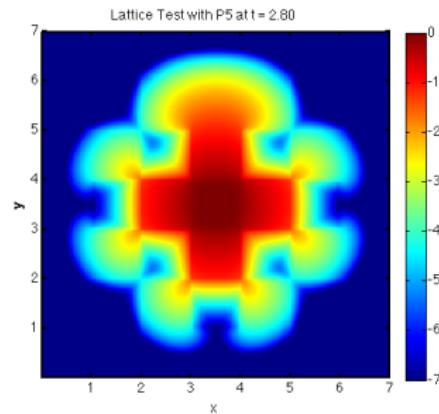
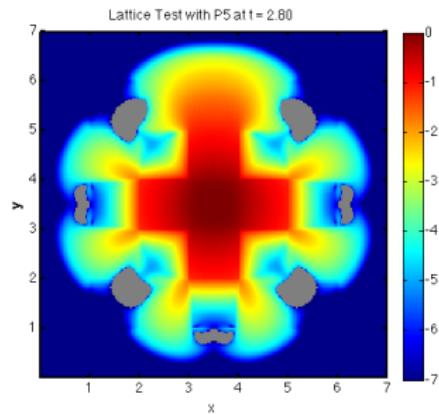
<sup>1</sup>McClarren, Hauck, JCP 2010

<sup>2</sup>Radice et al., JCP 2013

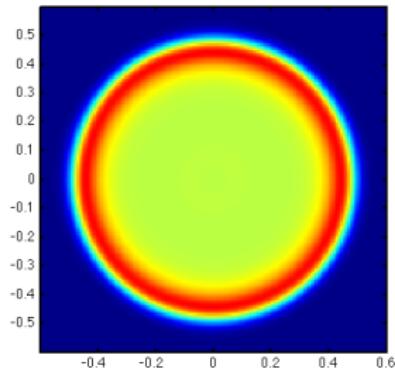
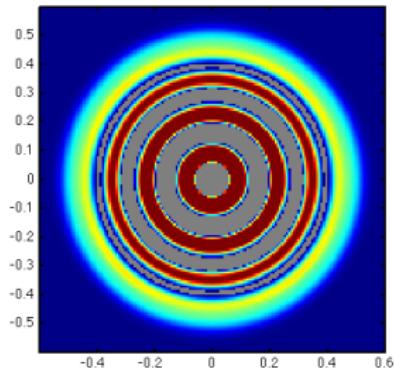
<sup>3</sup>Schmeiser, Zwirchmayr, SINUM 1999

<sup>4</sup>Seibold, Frank, TOMS 2014

# Checkerboard: $P_5$ versus $FP_5$



# Line Source: $P_9$ versus $FP_9$



# Challenges

## Challenges in radiation transport:

- Highly heterogeneous media
- Media/initial conditions/sources lead to non-smooth solutions
- Preserve realizability, rotational invariance
- Capture beams

## Challenges for spectral methods:

- Spectral methods achieve fast convergence for smooth solutions
- But suffer from the Gibbs phenomenon
- Idea of filtering: dampen the coefficients in the expansion
- **Con:** Some adjustments of the filter strength may be required for different problems
- **Pro:** Speed, overall accuracy, and simplicity

# FILTERED $P_N$

# Radiation Transport

$$\partial_t \psi(t, x, \Omega) + \Omega \cdot \nabla_x \psi(t, x, \Omega) + \sigma_a(x) \psi(t, x, \Omega) - (\mathcal{Q}\psi)(t, x, \Omega) = S(t, x, \Omega)$$

- $\psi(t, x, \Omega)$ : density of particles, with respect to the measure  $d\Omega dx$ , which at time  $t \in \mathbb{R}$  are located at position  $x \in \mathbb{R}^3$  and move in the direction  $\Omega \in \mathbb{S}^2$ .
- Scattering operator

$$(\mathcal{Q}\psi)(t, x, \Omega) = \sigma_s(x) \left[ \int_{\mathbb{S}^2} g(x, \Omega \cdot \Omega') \psi(t, x, \Omega') d\Omega' - \psi(t, x, \Omega) \right]$$

$$\mathcal{T}\psi = S$$

# Spherical Harmonic $P_N$ equations

## Notation:

- Real-valued spherical harmonic  $m_\ell^k$ ,  $\ell = 0, 1, \dots$ ,  
 $k = -\ell, \dots, \ell$
- Angular integration  $\langle \cdot \rangle = \int_{\mathbb{S}^2} (\cdot) d\Omega$

## Spectral Galerkin method:

- Expand unknown  $\psi \approx \psi_{PN} \equiv \mathbf{m}^T \mathbf{u}_{PN}$
- Plug into equation and project residual

$$\langle \mathbf{m}^T (\mathbf{m}^T \mathbf{u}_{PN}) \rangle = \langle \mathbf{m} S \rangle =: \mathbf{s}.$$

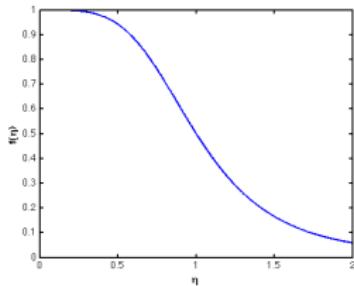
- Other combinations of ansatz and projection can be used!

## $P_N$ equations

$$\partial_t \mathbf{u}_{PN} + \mathbf{A} \cdot \nabla_x \mathbf{u}_{PN} + \sigma_a \mathbf{u}_{PN} - \sigma_s \mathbf{G} \mathbf{u}_{PN} = \mathbf{s},$$

where  $\mathbf{A} := \langle \mathbf{m} \mathbf{m}^T \Omega \rangle$  and  $\mathbf{G}$  is diagonal

# Filtering



- Filtering well-known in spectral methods
- A filter of order  $\alpha$  is a function  $f \in C^\alpha(\mathbb{R}^+)$ , which fulfills  $f(0) = 1$ ,  $f^{(k)}(0) = 0$ , for  $k = 1, \dots, \alpha - 1$ , and  $f^{(\alpha)}(0) \neq 0$
- Additional condition

$$f(\eta) \geq C(1 - \eta)^k, \quad \eta \in [\eta_0, 1]$$

- Filtering the expansion after every time step

$$\sum_{\ell=0}^N \sum_{k=-\ell}^{\ell} \left( f \left( \frac{\ell}{N+1} \right) \right)^{\beta \Delta t} u_\ell^k m_\ell^k.$$

# NUMERICAL ANALYSIS

# Main Result

## Galerkin estimate

$$\begin{aligned} \|\psi(t, \cdot, \cdot) - \psi_{\text{FPN}}(t, \cdot, \cdot)\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))} \\ \leq \|\psi(t, \cdot, \cdot) - \mathcal{P}_N \psi(t, \cdot, \cdot)\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))} \\ + t \left( \|\mathbf{a}_{N+1} \cdot \nabla_x \langle \mathbf{m}_{N+1} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \right. \\ \left. + \beta \|\mathbf{G}_f \langle \mathbf{m} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \right) \end{aligned}$$

## Rates

$$\begin{aligned} \|\psi(t, \cdot, \cdot) - \mathcal{P} \psi(t, \cdot, \cdot)\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))} &\leq C N^{-q} \|\psi\|_{C([0, T]; L^2(\mathbb{R}^3; H^q(\mathbb{S}^2)))} \\ \|\mathbf{a}_{N+1} \cdot \nabla_x \langle \mathbf{m}_{N+1} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \\ &\leq C N^{-r} \|\nabla_x \psi\|_{C([0, T]; L^2(\mathbb{R}^3; H^r(\mathbb{S}^2)))} \\ \|\mathbf{G}_f \langle \mathbf{m} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} &\leq \begin{cases} C N^{-q+1/2}, & \alpha > q - \frac{1}{2} \\ C N^{-\alpha+\varepsilon}, & \alpha \leq q - \frac{1}{2} \end{cases} \end{aligned}$$

# Sobolev Spaces

- $H^q(\mathbb{S}^2)$  Sobolev space on the unit sphere with norm

$$\|\Phi\|_{H^q(\mathbb{S}^2)} := \left( \sum_{|\alpha| \leq q} \int_{\mathbb{S}^2} |D^\alpha \Phi(\Omega)|^2 d\Omega \right)^{1/2}$$

- Spherical harmonics are eigenfunctions of Laplace-Beltrami operator

$$\mathcal{L}m_\ell^k = -\lambda_\ell m_\ell^k, \quad \lambda_\ell = \ell(\ell+1)$$

- Expansion coefficients  $\Phi_\ell^k := \langle m_\ell^k \Phi \rangle$  of any function  $\Phi \in H^{2q}(\mathbb{S}^2)$  satisfy

$$\Phi_\ell^k = \langle m_\ell^k \Phi \rangle = \frac{1}{(-\lambda_\ell)^q} \langle (\mathcal{L}^q m_\ell^k) \Phi \rangle = \frac{1}{(-\lambda_\ell)^q} \langle m_\ell^k \mathcal{L}^q \Phi \rangle$$

# Spectral Convergence

- $L^2$ -orthogonal projection of a generic function  $\Phi \in L^2(\mathbb{S}^2)$  onto  $\mathbb{P}_N$

$$\mathcal{P}_N \Phi = \mathbf{m}^T \langle \mathbf{m} \mathbf{m}^T \rangle^{-1} \langle \mathbf{m} \Phi \rangle = \mathbf{m}^T \langle \mathbf{m} \Phi \rangle$$

- Projection onto polynomials of exact degree  $\ell$   
 $(\mathcal{P}_\ell - \mathcal{P}_{\ell-1})\Phi = \mathbf{m}_\ell^T \langle \mathbf{m}_\ell \mathbf{m}_\ell^T \rangle^{-1} \langle \mathbf{m}_\ell \Phi \rangle = \mathbf{m}_\ell^T \langle \mathbf{m}_\ell \Phi \rangle$
- Spectral convergence

$$\begin{aligned}\|\langle \mathbf{m}_\ell \Phi \rangle\|_{\mathbb{R}^{n_\ell}}^2 &= \|(\mathcal{P}_\ell - \mathcal{P}_{\ell-1})\Phi\|_{L^2(\mathbb{S}^2)}^2 \leq \|(\mathcal{I} - \mathcal{P}_\ell)\Phi\|_{L^2(\mathbb{S}^2)}^2 \\ &= \sum_{k=\ell+1}^{\infty} |\phi_\ell|^2 = \sum_{k=\ell+1}^{\infty} \frac{1}{(-\lambda_\ell)^{2q}} |\langle \mathbf{m}_\ell \mathcal{L}^q \Phi \rangle|^2 \\ &\leq \frac{1}{(\ell(\ell+1))^{2q}} \|\phi\|_{H^{2q}(\mathbb{S}^2)}^2\end{aligned}$$

## Step 1: Modified Equation

- Time step

$$\mathbf{u}_{\text{FPN}}^{n+1,*} = \mathbf{u}_{\text{FPN}}^n - \Delta t (\mathbf{A} \cdot \nabla_x \mathbf{u}_{\text{FPN}}^n + \sigma_a \mathbf{u}_{\text{FPN}}^n - \sigma_s \mathbf{G} \mathbf{u}_{\text{FPN}}^n - \mathbf{s}^n)$$

- Filtering

$$\mathbf{u}_{\text{FPN}}^{n+1} = \mathbf{f}^{\beta \Delta t} \mathbf{u}_{\text{FPN}}^{n+1,*} = \mathbf{u}_{\text{FPN}}^{n+1,*} + \Delta t \frac{\exp(\beta \log(\mathbf{f}) \Delta t) - 1}{\Delta t} \mathbf{u}_{\text{FPN}}^{n+1,*}$$

- Operator split discretization of

### Modified equation

$$\partial_t \mathbf{u}_{\text{FPN}} + \mathbf{A} \cdot \nabla_x \mathbf{u}_{\text{FPN}} + \sigma_a \mathbf{u}_{\text{FPN}} - \sigma_s \mathbf{G} \mathbf{u}_{\text{FPN}} - \beta \mathbf{G}_f \mathbf{u}_{\text{FPN}} = \mathbf{s},$$

where  $\mathbf{G}_f$  is diagonal with entries  $\log\left(f\left(\frac{\ell}{N+1}\right)\right)$ ,  $\ell = 0, \dots, N$ .

## Step 2: Galerkin Estimate

- Residual

$$\psi - \psi_{\text{FPN}} = (\psi - \mathcal{P}_N \psi) + \mathcal{P}_N \psi - \psi_{\text{FPN}} = (\psi - \mathcal{P}_N \psi) + \mathbf{m}^T \mathbf{r}$$

- Multiply by  $\mathbf{m}^T \mathbf{r}$  and integrate in angle and space

$$\begin{aligned} \frac{1}{2} \partial_t \int_{\mathbb{R}^3} |\mathbf{r}|^2 dx &= - \int_{\mathbb{R}^3} \mathbf{r}_N^T \mathbf{a}_{N+1} \cdot \nabla_x \langle \mathbf{m}_{N+1} \psi \rangle dx \\ &\quad - \sigma_f \int_{\mathbb{R}^3} \mathbf{r}^T \mathbf{G}_f \langle \mathbf{m} \psi \rangle dx - \int_{\mathbb{R}^3} \mathbf{r}^T \mathbf{M} \mathbf{r} dx . \end{aligned}$$

- $\mathbf{M} := \sigma_a \mathbf{I} - \sigma_s \mathbf{G} - \sigma_f \mathbf{G}_f$  is positive definite

- This yields

$$\begin{aligned} \partial_t \|\mathbf{r}\|_{L^2(\mathbb{R}^3; \mathbb{R}^n)} &\leq \|\mathbf{a}_{N+1} \cdot \nabla_x \langle \mathbf{m}_{N+1} \psi \rangle\|_{L^2(\mathbb{R}^3; \mathbb{R}^{2N+1})} \\ &\quad + \sigma_f \|\mathbf{G}_f \langle \mathbf{m} \psi \rangle\|_{L^2(\mathbb{R}^3; \mathbb{R}^n)} \end{aligned}$$

- Control error by projection error + residual  $\mathbf{r}$

## Step 3: Convergence Estimate

- Estimate filter term

$$\begin{aligned} & \| \mathbf{G}_f \langle \mathbf{m} \psi(t, \cdot, \cdot) \rangle \|_{L^2(\mathbb{R}^3; \mathbb{R}^n)}^2 \\ &= \sum_{\ell=0}^N \log^2 \left( f \left( \frac{\ell}{N+1} \right) \right) \| \langle \mathbf{m}_\ell \psi(t, \cdot, \cdot) \rangle \|_{L^2(\mathbb{R}^3; \mathbb{R}^{n_\ell})}^2 \\ &= \sum_{\ell=1}^N \log^2 \left( f \left( \frac{\ell}{N+1} \right) \right) \| (\mathcal{P}_\ell - \mathcal{P}_{\ell-1}) \psi(t, \cdot, \cdot) \|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))}^2 \\ &= C \sum_{\ell=1}^N \log^2 \left( f \left( \frac{\ell}{N+1} \right) \right) \| (\mathcal{I} - \mathcal{P}_{\ell-1}) \psi(t, \cdot, \cdot) \|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))}^2 \\ &\leq C \sum_{\ell=1}^N \log^2 \left( f \left( \frac{\ell}{N+1} \right) \right) \frac{1}{\ell^{2q}} \| \psi(t, \cdot, \cdot) \|_{L^2(\mathbb{R}^3; H^q(\mathbb{S}^2))}^2 \end{aligned}$$

## Step 3: Convergence Estimate

- For  $\theta \leq 2q$

$$\begin{aligned} & \sum_{\ell=1}^N \log^2 \left( f\left(\frac{\ell}{N+1}\right) \right) \frac{1}{\ell^{2q}} \\ & \leq \frac{1}{(N+1)^{\theta-1}} \underbrace{\frac{1}{N+1} \sum_{\ell=1}^N \log^2 \left( f\left(\frac{\ell}{N+1}\right) \right) \left(\frac{N+1}{\ell}\right)^\theta}_{=: \Sigma} \end{aligned}$$

- Interpret as Riemann sum

$$\Sigma \sim \int_0^1 \log^2(f(\eta)) \eta^{-\theta} d\eta$$

- Around  $\eta = 0$ ,  $\log f(\eta) \leq C\eta^\alpha$
- $\Sigma$  Integrable for  $\theta < 2\alpha + 1$

## Step 3: Convergence Estimate

Two cases:

Case 1:  $\alpha > q - \frac{1}{2}$ . Choose  $\theta = 2q$ , convergence limited by the regularity of  $\psi$

$$\|\mathbf{G}_f \langle \mathbf{m} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \leq CN^{-q+1/2}$$

Case 2:  $\alpha \leq q - \frac{1}{2}$ . Choose  $\theta = 2\alpha + 1 - \delta$ , where  $\delta > 0$  is arbitrary, convergence limited by the filter order

$$\|\mathbf{G}_f \langle \mathbf{m} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \leq CN^{-\alpha+\varepsilon},$$

where  $\varepsilon = \delta/2$

# Main Result

## Galerkin estimate

$$\begin{aligned} \|\psi(t, \cdot, \cdot) - \psi_{\text{FPN}}(t, \cdot, \cdot)\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))} \\ \leq \|\psi(t, \cdot, \cdot) - \mathcal{P}\psi(t, \cdot, \cdot)\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))} \\ + t(\|\mathbf{a}_{N+1} \cdot \nabla_x \langle \mathbf{m}_{N+1} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \\ + \beta \|\mathbf{G}_f \langle \mathbf{m} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))}) , \end{aligned}$$

## Rates

$$\begin{aligned} \|\psi(t, \cdot, \cdot) - \mathcal{P}\psi(t, \cdot, \cdot)\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))} &\leq CN^{-q} \|\psi\|_{C([0, T]; L^2(\mathbb{R}^3; H^q(\mathbb{S}^2)))} \\ \|\mathbf{a}_{N+1} \cdot \nabla_x \langle \mathbf{m}_{N+1} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \\ &\leq CN^{-r} \|\nabla_x \psi\|_{C([0, T]; L^2(\mathbb{R}^3; H^r(\mathbb{S}^2)))} \\ \|\mathbf{G}_f \langle \mathbf{m} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} &\leq \begin{cases} CN^{-q+1/2}, & \alpha > q - \frac{1}{2} \\ CN^{-\alpha+\varepsilon}, & \alpha \leq q - \frac{1}{2} \end{cases} \end{aligned}$$

# Sharper Estimate

## Galerkin estimate

$$\begin{aligned} \|\psi(t, \cdot, \cdot) - \psi_{\text{FPN}}(t, \cdot, \cdot)\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))} \\ \leq \|\psi(t, \cdot, \cdot) - \mathcal{P}\psi(t, \cdot, \cdot)\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))} \\ + t \left( \|\mathbf{a}_{N+1} \cdot \nabla_x \langle \mathbf{m}_{N+1} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \right. \\ \left. + \beta \|\mathbf{G}_f \langle \mathbf{m} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \right), \end{aligned}$$

## Rates for monotone moment sequences

$$\begin{aligned} \|\psi(t, \cdot, \cdot) - \mathcal{P}\psi(t, \cdot, \cdot)\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))} &\leq CN^{-q} \|\psi\|_{C([0, T]; L^2(\mathbb{R}^3; H^q(\mathbb{S}^2)))} \\ \|\mathbf{a}_{N+1} \cdot \nabla_x \langle \mathbf{m}_{N+1} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} \\ &\leq C N^{-(r+\frac{1}{2})} \|\nabla_x \psi\|_{C([0, T]; L^2(\mathbb{R}^3; H^r(\mathbb{S}^2)))} \\ \|\mathbf{G}_f \langle \mathbf{m} \psi \rangle\|_{C([0, T]; L^2(\mathbb{R}^3; \mathbb{R}^n))} &\leq \begin{cases} C N^{-q}, & \alpha > q \\ CN^{-\alpha+\varepsilon}, & \alpha \leq q \end{cases} \end{aligned}$$

# NUMERICAL RESULTS

## Numerical Results: General setup

- Use the code StaRMAP to compute the  $P_N$  and  $FP_N$  solutions for  $N = 3, 5, 17, 33, (65)$ .
- Apply the filter term after each sub-step to the updated components.
- Use the exponential filter of order  $\alpha = 2, 4, 8, 16$   
$$f(\eta) = \exp(c\eta^\alpha), \text{ with } c = \log(\varepsilon_M)$$

with  $\varepsilon_M$  being the machine precision. Set the effective filter opacity  $f_{\text{eff}} = 10$  ( $f_{\text{eff}} = \beta \log(f(\frac{N}{N+1}))$ ).

- Fix the spatial resolution, so that the space-time errors are negligibly small
- Compare to reference solution  $P_{N_{\text{true}}}$
- Highest resolution  $P_{129}$  (8515 moments) on  $500 \times 500$  grid (altogether  $2.1 \times 10^9$  unknowns)

# Numerical Results: Estimates

- Measure smoothness of true solution

$$B_N = \|\langle \mathbf{m}_N \psi \rangle\|_{L(\mathbb{R}^2, \mathbb{R}^n)} \sim N^{-q+\frac{1}{2}}$$

$$D_N = \|\langle \mathbf{m}_N \nabla_x \psi \rangle\|_{L(\mathbb{R}^2, \mathbb{R}^n)} \sim N^{-r+\frac{1}{2}}$$

- Compare to convergence estimate

$$E_N = \|\psi - \psi_N\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))}$$

$$R_N = \|\mathcal{P}\psi - \psi_N\|_{L^2(\mathbb{R}^3; L^2(\mathbb{S}^2))}$$

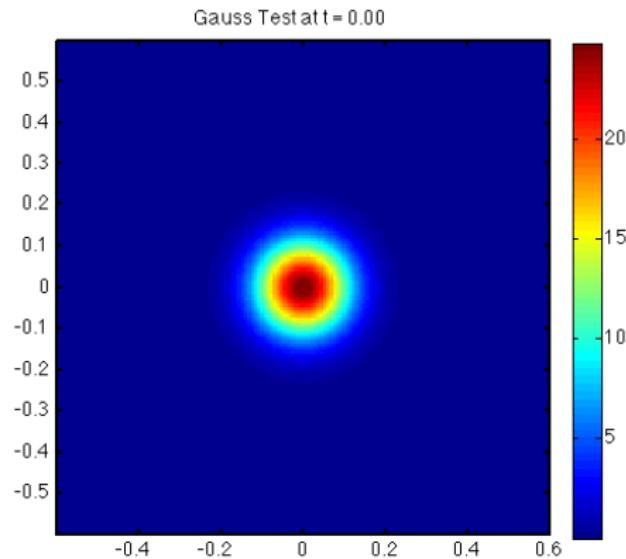
- Expectation

With filter:  $E_N \sim R_N \sim N^{-\min\{q, r + \frac{1}{2}, \alpha\}}$

Without filter:  $E_N \sim N^{-\min\{q, r + \frac{1}{2}\}}, R_N \sim N^{-(r + \frac{1}{2})}$

- Central difference for  $\nabla_x$ , trapezoidal rule for integration

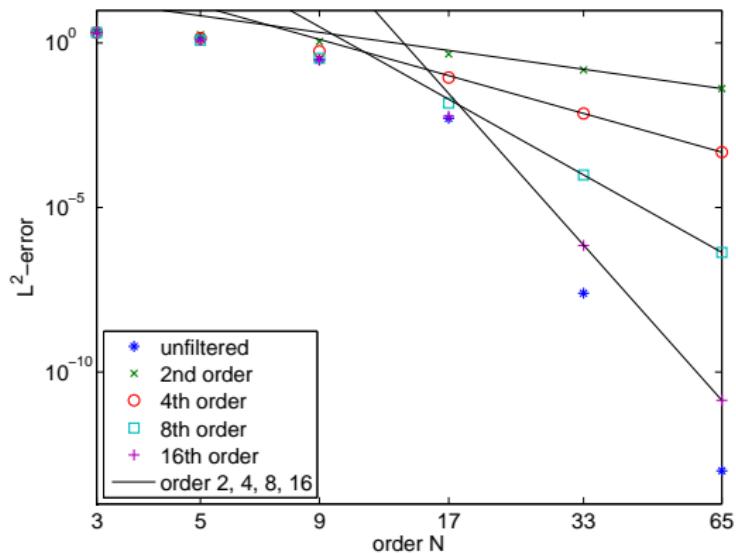
# Gaussian Test: Setup



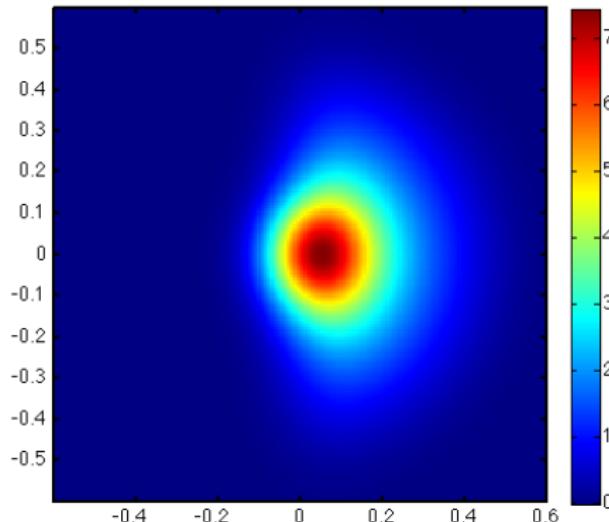
- Initial condition:  $u_0^0 = \frac{1}{4\pi \times 10^{-3}} \exp\left(-\frac{x^2+y^2}{4 \times 10^{-3}}\right)$ ,  
 $u_\ell^k = 0$ , for  $k, \ell \neq 0$
- Purely scattering medium:  $\sigma_t = \sigma_s = 1$

# Gaussian Test: Results

- $q = r = \infty$
- With filter:  $E_N \sim R_N \sim N^{-\min\{q, r + \frac{1}{2}, \alpha\}}$
- Without filter:  $E_N \sim N^{-\min\{q, r + \frac{1}{2}\}}, R_N \sim N^{-(r + \frac{1}{2})}$

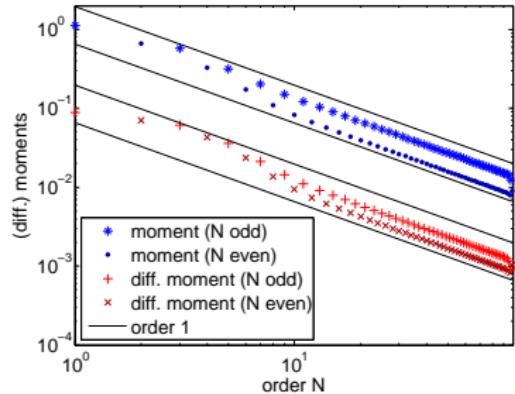


# Hemisphere Test: Setup

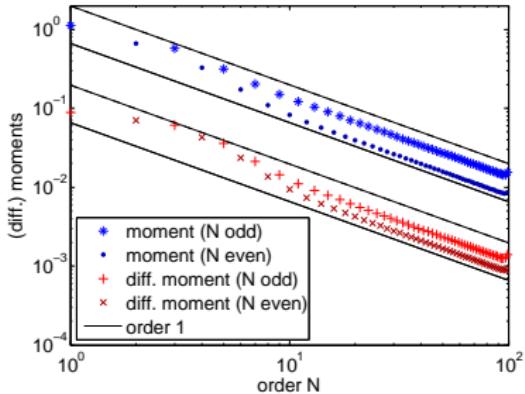


- Source term:  $S(t, x, \Omega) = \frac{1}{4\pi \times 10^{-3}} \exp\left(-\frac{x^2+y^2}{4 \times 10^{-3}}\right) \chi_{\mathbb{R}^+}(\Omega_x)$
- Vacuum:  $\sigma_t = 0$ .

# Hemisphere Test: Smoothness



(a)  $P_{98}$



(b)  $P_{99}$

- From  $B_N \sim N^{-q+\frac{1}{2}}$  and  $D_N \sim N^{-r+\frac{1}{2}}$  we conclude  
$$q \approx r \approx 0.5$$

# Hemisphere Test: Results

Filter order	$\mathcal{E}_3^5$	$\mathcal{E}_5^9$	$\mathcal{E}_9^{17}$	$\mathcal{E}_{17}^{33}$	$\mathcal{R}_3^5$	$\mathcal{R}_5^9$	$\mathcal{R}_9^{17}$	$\mathcal{R}_{17}^{33}$
2	0.55	0.58	0.57	0.58	0.44	0.61	0.59	0.52
4	0.67	0.60	0.55	0.61	0.71	0.70	0.57	0.52
8	0.75	0.61	0.56	0.63	1.06	0.83	0.61	0.56
16	0.77	0.64	0.57	0.64	1.14	1.03	0.79	0.64
$\infty$	0.71	0.59	0.56	0.65	1.33	1.26	0.99	0.96

- $q \approx r \approx 0.5$
- With filter:

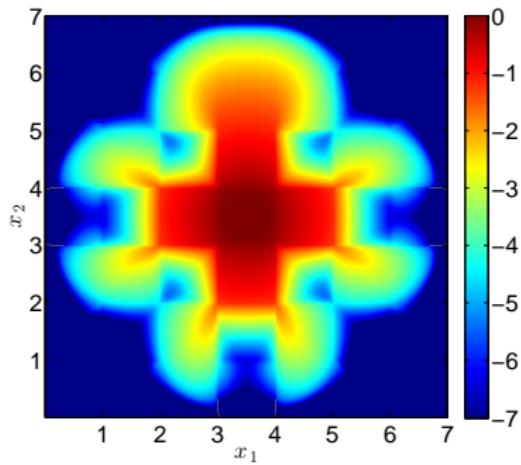
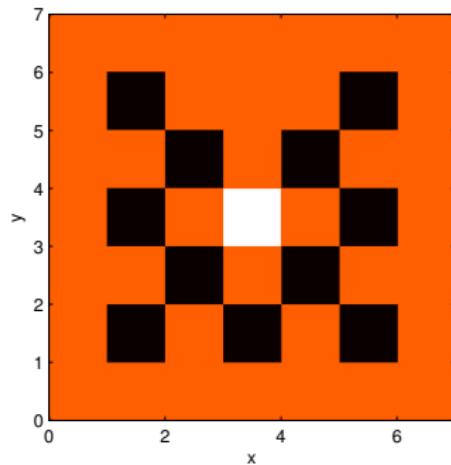
$$E_N \sim R_N \sim N^{-\min\{q, r + \frac{1}{2}, \alpha\}}$$

- Without filter:

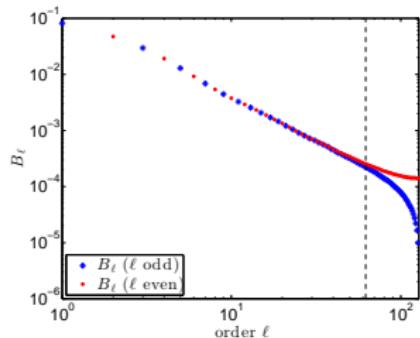
$$E_N \sim N^{-\min\{q, r + \frac{1}{2}\}}$$

$$R_N \sim N^{-(r + \frac{1}{2})}$$

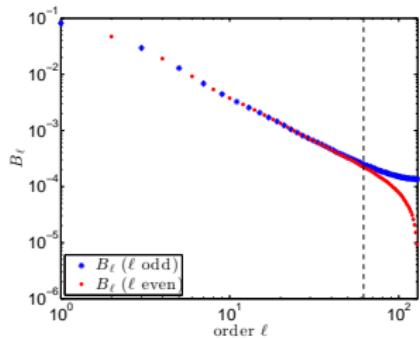
# Checkerboard Test: Setup



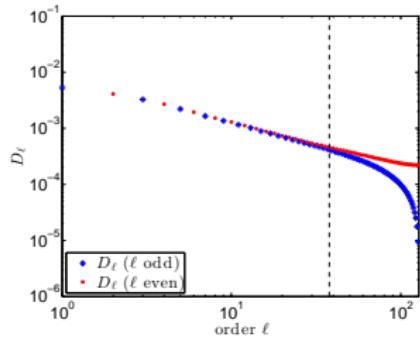
# Checkerboard Test: Smoothness



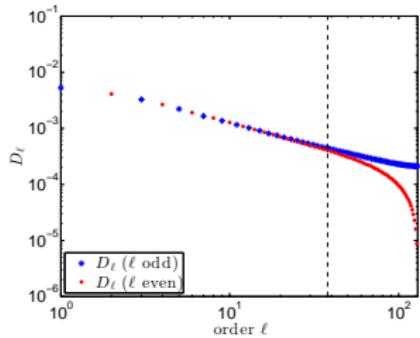
(c)  $B_\ell$  computed with  $P_{128}$



(d)  $B_\ell$  computed with  $P_{129}$



(e)  $D_\ell$  computed with  $P_{128}$



(f)  $D_\ell$  computed with  $P_{129}$

## Checkerboard Test: More Smoothness

$(N_1, N_2)$	$B_{N_1}^{N_2}$	$D_{N_1}^{N_2}$
(2,4)	1.3188	0.6213
(4,8)	1.8212	0.8161
(8,16)	1.5208	0.8293
(16,32)	1.5782	0.8679

(a) even order moments

$(N_1, N_2)$	$B_{N_1}^{N_2}$	$D_{N_1}^{N_2}$
(3,5)	1.6167	0.7818
(5,9)	1.8371	0.8204
(9,17)	1.4901	0.7998
(17,33)	1.5511	0.7691

(b) odd order moments

- From  $B_N \sim N^{-q+\frac{1}{2}}$  and  $D_N \sim N^{-r+\frac{1}{2}}$  we conclude  
 $q \approx 1.0$  and  $r \approx 0.25$

## Checkerboard Test: Results

Filter order	$\mathcal{E}_3^5$	$\mathcal{E}_5^9$	$\mathcal{E}_9^{17}$	$\mathcal{E}_{17}^{33}$	$\mathcal{R}_3^5$	$\mathcal{R}_5^9$	$\mathcal{R}_9^{17}$	$\mathcal{R}_{17}^{33}$
2	0.89	0.80	0.94	1.05	0.86	0.78	0.93	1.05
4	1.02	1.15	1.13	1.05	0.98	1.21	1.21	1.06
8	1.20	1.22	1.04	1.06	1.32	1.55	1.14	1.16
16	1.61	1.31	1.03	1.04	2.10	2.12	1.23	1.20
$\infty$	1.10	0.95	0.98	1.00	1.10	0.85	0.95	0.96

- $q = 1, r = 0.25$
- With filter:

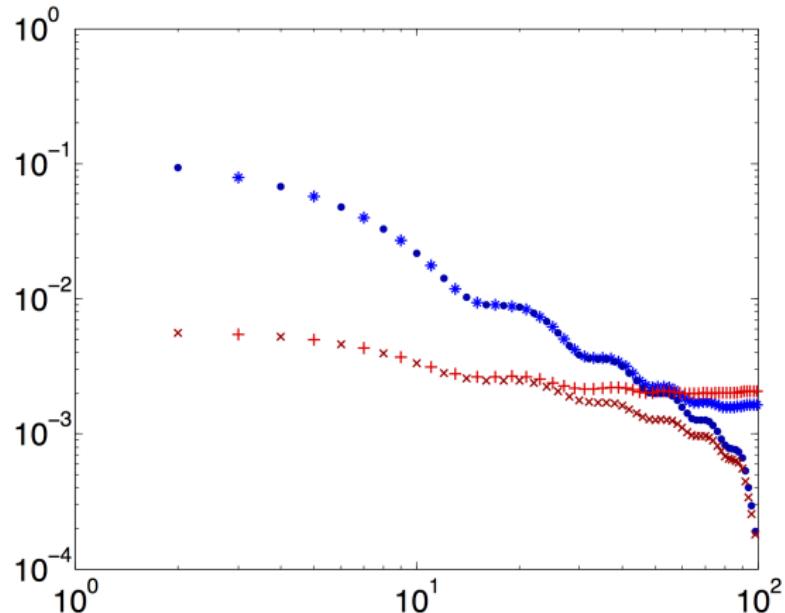
$$E_N \sim R_N \sim N^{-\min\{q, r + \frac{1}{2}, \alpha\}}$$

- Without filter:

$$E_N \sim N^{-\min\{q, r + \frac{1}{2}\}}$$

$$R_N \sim N^{-(r + \frac{1}{2})}$$

# Things Aren't Always So Clear



Box source instead of Gaussian.

# Summary & Outlook

## Summary:

- Proof of global  $L^2$  convergence rates for filtered spherical harmonic ( $FP_N$ ) equations
- Dependence of the convergence rates on regularity of transport solution and order of the filter
- Highly resolved numerical experiments are pretty much in agreement with theoretical predictions

## Outlook:

- Local analysis to show improvements by filtering
- Similar analysis for entropy or other non-linear closures